

Numerical study of the performance of a church window tube bundle condenser

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Abstract—Two steam surface condenser design methods are compared: that of the Heat Exchange Institute Standards and a three-dimensional steady-state numerical model for simulating the flow, the heat and the mass transfer in this type of industrial equipment. In the numerical model, the study of the shell flow is carried out using a single-phase approach, the presence of the tubes is simulated through the porosity factor, and the heat and mass transfer processes are modelled using empirical correlations. The variables condenser pressure and overall heat transfer coefficient are analysed applying both methods. The influence on these variables of the steam mass flow rate from the turbine exhaust, the venting outlet pressure, the tube cleanliness factor, the inlet water temperature, and the water velocity is studied. The overall heat transfer coefficient is under-predicted with the numerical model, its value being 4.5–8.5% lower than that obtained using the Heat Exchange Institute Standards method. © 2001 Éditions scientifiques et médicales Elsevier SAS

condenser performance / steam surface condenser / numerical modelling / condensation / noncondensable gases

Nomenclature

A	total heat transfer area	m^2
C_p	specific heat at constant pressure . .	$\text{J}\cdot\text{kg}^{-1}\cdot\text{K}^{-1}$
D_i	inner diameter of tubes	m
D_o	outer diameter of tubes	m
D_{sa}	diffusivity of steam in air	$\text{m}^2\cdot\text{s}^{-1}$
F	flow resistance force per unit volume due to the tube bundle	$\text{N}\cdot\text{m}^{-3}$
F_C	tube cleanliness factor	
F_M	correction factor for material and tube gauge	
F_W	correction factor for inlet water temperature	
H	Ackerman's correction coefficient	
h	heat transfer coefficient	$\text{W}\cdot\text{m}^{-2}\cdot\text{K}^{-1}$
i	specific enthalpy	$\text{J}\cdot\text{kg}^{-1}$
k	thermal conductivity of tubes	$\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$
LMTD	logarithmic mean temperature difference	K
m	mass flow rate	$\text{kg}\cdot\text{s}^{-1}$
m_a	mass flow rate per unit area	$\text{kg}\cdot\text{s}^{-1}\cdot\text{m}^{-2}$

m_v	mass flow rate per unit volume	$\text{kg}\cdot\text{s}^{-1}\cdot\text{m}^{-3}$
N	total number of control volumes	
P	pressure	Pa
P_t	tube pitch	m
Q	heat transfer rate	W
q_a	heat transfer rate per unit area	$\text{W}\cdot\text{m}^{-2}$
q_v	heat transfer rate per unit volume	$\text{W}\cdot\text{m}^{-3}$
R	thermal resistance $= 1/U$	$\text{m}^2\cdot\text{K}\cdot\text{W}^{-1}$
R_f	fouling resistance	$\text{m}^2\cdot\text{K}\cdot\text{W}^{-1}$
T	temperature	$^{\circ}\text{C}$
\mathbf{u}	velocity vector	$\text{m}\cdot\text{s}^{-1}$
u	velocity	$\text{m}\cdot\text{s}^{-1}$
U	overall heat transfer coefficient	$\text{W}\cdot\text{m}^{-2}\cdot\text{K}^{-1}$
U_1	uncorrected overall heat transfer coefficient	$\text{W}\cdot\text{m}^{-2}\cdot\text{K}^{-1}$
V	volume	m^3
Y_a	air mass fraction	

Greek symbols

β	volumetric porosity factor	
ε_P	relative error of pressure $= (P_{NM} - P_{SSSC})/P_{SSSC} \cdot 100$	%
ε_U	relative error of overall heat transfer coefficient $= (U_{NM} - U_{SSSC})/U_{SSSC} \cdot 100$	%
ρ	density	$\text{kg}\cdot\text{m}^{-3}$

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μ_{eff}	effective dynamic viscosity of the mixture = $\mu_L + \mu_T$	$\text{N}\cdot\text{m}^{-2}\cdot\text{s}^{-1}$
μ_L	laminar dynamic viscosity of the mixture	$\text{N}\cdot\text{m}^{-2}\cdot\text{s}^{-1}$
μ_T	turbulent dynamic viscosity of the mixture = $20 \mu_L$	$\text{N}\cdot\text{m}^{-2}\cdot\text{s}^{-1}$

Subscripts

avg	average value
c	condensate
cc	corrected condensate
f	fouling
g	gas (steam–air mixture)
i	inlet
if	mixture–condensate interface
j	index
lg	liquid–gas change
NM	numerical model
o	outlet
s	steam
SSSC	“Standards for Steam Surface Condensers”
w	water
x, y, z	coordinates

1. INTRODUCTION

The design of thermal power plant condensers has traditionally been based on recommendations elaborated by condenser manufacturers founded on previous designs, and expensive experimental tests. In “Standards for Steam Surface Condensers” (SSSC) developed by the Heat Exchange Institute (HEI) [1], the following equations (1)–(4) are proposed for condenser design. The heat transfer rate in the condenser, the saturation pressure, the water mass flow rate and the inlet water temperature are required to solve these equations. The condensate mass flow rate, the outlet water temperature, and the total heat transfer area are calculated, together with the definition of the arrangement, the material and the geometric parameters of the tubes:

$$Q = m_c i_{lg} \quad (1)$$

$$Q = m_w C_{p_w} (T_{w,o} - T_{w,i}) \quad (2)$$

$$Q = U A \frac{T_{w,o} - T_{w,i}}{\ln((T_g - T_{w,i})/(T_g - T_{w,o}))} \quad (3)$$

$$U = U_1 F_W F_M F_C \quad (4)$$

where Q is the heat transfer rate in the condenser, m_c is the condensate mass flow rate, i_{lg} is the steam vapourisation specific enthalpy, m_w is the water mass

flow rate, C_{p_w} is the specific heat at constant pressure of the water evaluated at the average water temperature between the inlet and the outlet, $T_{w,o}$ is the outlet water temperature, $T_{w,i}$ is the inlet water temperature, U is the overall heat transfer coefficient, A is the total heat transfer area, T_g is the mixture temperature calculated as the saturation temperature corresponding to the steam pressure, U_1 is the uncorrected overall heat transfer coefficient, F_W is the correction factor for the inlet water temperature, F_M is the correction factor for the material and tube gauge and F_C is the tube cleanliness factor.

The SSSC does not take into account a number of factors that affect the condenser pressure and the overall heat transfer coefficient: the presence of noncondensable gases mixed with the steam, inundation of the tubes by the condensing steam, the pressure loss of the mixture across the tube bundle, the geometry of the tube bundle, and the design of the noncondensable gases extraction equipment.

There is a considerable number of numerical models for simulating the performance of steam surface condensers [2–9], which give detailed information on the distribution of the mixture velocity, the mixture pressure, the concentration of noncondensable gases, and the water temperature. These numerical models represent a very useful tool for studying the processes inside the condensers, and for investigating the effect of alternative geometries more quickly and cheaply than experimental studies, with a reduced number of tests. Most of the models assume two-dimensional performance, and use a single-phase approach to study the flow within the shell. The thermal effects of steam condensation, heat and mass transfer from the mixture to the condensate at the interface, steam shear at the interface, inundation of the tubes by the condensing steam, and pressure loss of the mixture across the tube bundle are modelled through empirical correlations. The effect of tube fouling is included as an additional thermal resistance to the heat transfer, which could be quite substantial depending on the quality of the water, particularly in condensers working with seawater as coolant. The performance of this type of model has been experimentally assessed in [2, 8, 9].

In the present study, a numerical model is described that is based on equivalent assumptions to those mentioned above, but which considers the flow inside the shell as three dimensional, using an approach based on the Colburn–Hougen model [10] and empirical correlations for convective heat transfer coefficients. The model was applied to a feedwater heater from a thermal power plant, and the numerical results were found to be satisfactory compared with design and experimental data in [11].

The focus of the study is now to compare the results of the SSSC with those obtained using the numerical model (NM), since there are no existing bibliographic references in this field. Both methods are applied to a single-pass condenser with a church window tube bundle. The comparison is based on the analysis of two variables: the condenser pressure and the overall heat transfer coefficient. The evolution of the variables is studied with respect to five parameters: the steam mass flow rate from the turbine exhaust, the venting outlet pressure, the tube cleanliness factor, the inlet water temperature, and the water velocity. Moreover, the differences between the values of the variables obtained using both methods are quantified.

2. NUMERICAL MODEL

2.1. Description

The numerical model simulates the steady-state flow and the heat and mass transfer processes in thermal power plant condensers.

The shell-fluid is a mixture of steam and noncondensable gases, mainly air. The flow within the shell is simulated using a single-phase approach, assuming that the condensate disappears from the domain as it is formed.

The tube-fluid is cooling water. As the flow pattern is known, the water velocity along the tubes is considered constant, and it is not necessary to obtain the water velocity field.

The conservation equations of the mixture mass (5), the mixture momentum (6), the air mass (7), and the water energy (8) are solved in each control volume:

$$\text{div}(\beta \rho_g \mathbf{u}_g) = -\beta m_{c,v} \quad (5)$$

$$\text{div}(\beta \rho_g (\mathbf{u}_g \otimes \mathbf{u}_g)) = \text{div}(\beta \mu_{\text{eff}} \text{grad}(\mathbf{u}_g)) - \beta \text{grad}(P_g) - \beta \mathbf{F} - \beta m_{c,v} \mathbf{u}_g \quad (6)$$

$$\text{div}(\beta \rho_g \mathbf{u}_g Y_a) = \text{div}(\beta \rho_g D_{sa} \text{grad}(Y_a)) \quad (7)$$

$$\text{div}(\rho_w \mathbf{u}_w i_w) = q_v \quad (8)$$

where β is the porosity factor, ρ_g is the mixture density, \mathbf{u}_g is the mixture velocity, $m_{c,v}$ is the condensate mass flow rate per unit volume, μ_{eff} is the effective dynamic viscosity of the mixture, P_g is the mixture pressure, \mathbf{F} is the flow resistance force due to the tube bundle, Y_a is the air mass fraction, D_{sa} is the diffusivity of steam

in air, ρ_w is the water density, \mathbf{u}_w is the water velocity, and q_v is the heat transfer rate per unit volume.

The tube bundle is simulated using the volumetric porosity factor, which is defined as the ratio of the volume occupied by the shell-side fluid and the total volume. It is given by equation (9) for a staggered arrangement of tubes at 60°:

$$\beta = 1 - \frac{\pi}{2\sqrt{3}} \left(\frac{D_o}{Pt} \right)^2 \quad (9)$$

where β is the volumetric porosity factor, D_o is the outer diameter of the tubes, and Pt is the tube pitch.

The turbulence within the shell is modelled through the effective dynamic viscosity, which is defined as the sum of the laminar and turbulent dynamic viscosities. In the present study, the turbulent dynamic viscosity is defined as 20 times the laminar dynamic viscosity, in accordance with the sensitivity analysis carried out in prior studies [8, 9, 12].

The condensate mass flow rate is taken into account in the conservation equations of the mixture mass and momentum as source terms, which represent mass and momentum sinks.

The friction force in the source term of the momentum conservation equation is modelled using adequate empirical correlations for crossflow and longitudinal flow as a function of geometric parameters of the tubes, friction factors, and physical properties of the mixture. The expressions in [5, 8, 9] are used for calculating the cross-flow components, and the internal flow expression is employed for evaluating the longitudinal component.

The study of the heat and mass transfer is carried out according to Colburn and Hougen's film theory [12], taking into account the thermal resistance of the water inside the tubes, the wall of the tubes, the condensate film, tube fouling, and the gaseous mixture, as well as the vaporisation heat transfer rate. Equation (10) is solved in each control volume to obtain the temperature at the mixture–condensate interface:

$$\begin{aligned} q_a &= \frac{T_{if} - T_w}{\frac{D_o}{D_i} \frac{1}{h_w} + \frac{D_o}{2k} \ln \left(\frac{D_o}{D_i} \right) + \frac{1}{h_{cc}} + R_{f,NM}} \\ &= h_g \frac{H}{1 - e^{-H}} (T_g - T_{if}) + i_{lg} m_{c,a} \end{aligned} \quad (10)$$

where q_a is the heat transfer rate per unit area, T_{if} is the temperature at the mixture–condensate interface, T_w is the water temperature, D_i is the inner diameter of the tubes, h_w is the water-side heat transfer coefficient, k is the thermal conductivity of the tubes, h_{cc} is the

condensate film heat transfer coefficient corrected by the condensate inundation of the tubes factor and the steam shear at the interface factor, $R_{f,NM}$ is the tube fouling thermal resistance, h_g is the heat transfer coefficient of the mixture, H is Ackerman's correction factor, i_{lg} is the vaporisation specific enthalpy of steam, and $m_{c,a}$ is the condensate mass flow rate per unit area.

The water-side heat transfer coefficient is calculated using the Dittus–Boelter correlation given in [13].

The condensate film heat transfer coefficient is calculated from Nusselt's correlation [14], and is corrected using two factors to take into account the effects of the condensate inundation of the tubes [15], and the steam shear at the interface [14].

The heat transfer coefficient of the mixture is evaluated from the correlation proposed by Taborek [16], and is corrected with Ackerman's factor [17] to take into account the diffusion effects due to the presence of air mixed with the steam.

2.2. Application

The three-dimensional numerical model is applied to a condenser with a horizontal tube bundle. *Figure 1(a)* shows a cross section of the condenser. The steam from the turbine exhaust enters the condenser at the top of the shell and flows across the tubes to the bundle centre area due to the vacuum produced by the venting equipment, which extracts the noncondensable gases and the noncondensed steam. The mixture passes to the venting conduct through a set of orifices of very small diameter uniformly distributed along the conduct. The aluminum brass tubes are in a staggered arrangement at 60°, 9.75 m long, 0.024 m outer diameter, and 18 BWG gauge.

The geometry and the boundary conditions are symmetrical with respect to the y-direction and only one half is simulated. The Cartesian grid used is 0.1 m in size, except in the baffles and venting area, where a refined grid is defined to approximate the baffle shape, and to take into account the small dimension of the venting area. The resulting grid is $23 \times 87 \times 120$ control volumes, and a tube bundle cross section is shown in *figure 1(b)*.

The boundary conditions defined are the following: the steam and air mass flow rates in the shell inlet, the water velocity and temperature in the tube inlet, and the venting outlet pressure.

The operating conditions are indicated in *table I*. The air mass flow rate in the condenser is assumed to

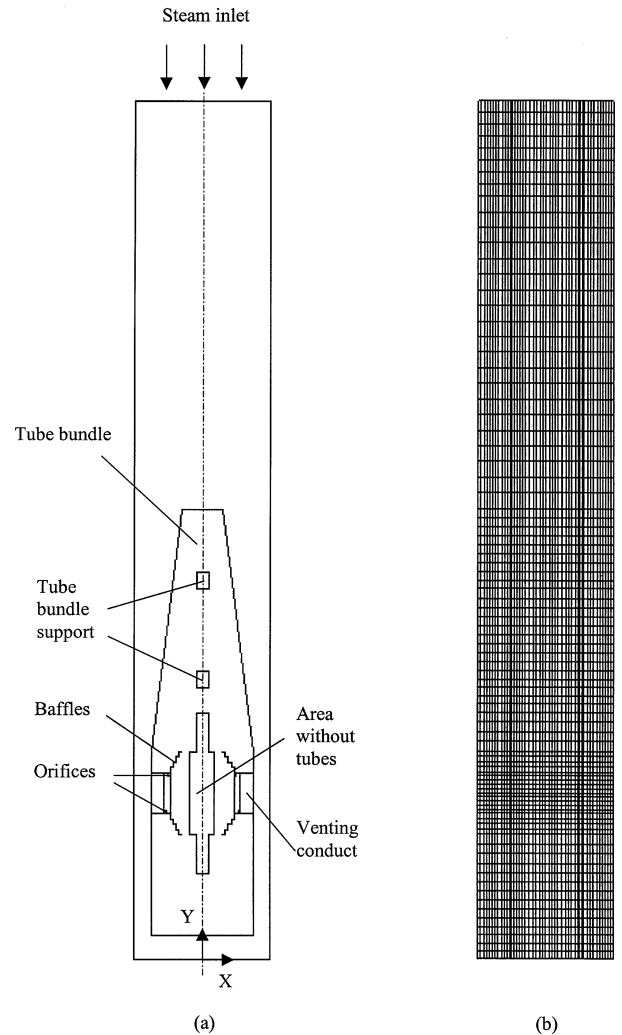


Figure 1. (a) Geometry of the condenser, (b) grid.

TABLE I
Performance conditions.

Constant data			
Air mass flow rate (kg·s ⁻¹)	5.56·10 ⁻⁴		
Variable data			
Steam mass flow rate (kg·s ⁻¹)	30.84	23.73	16.95
Venting outlet pressure (Pa)	2800	3200	3800
Tube cleanliness factor (dimensionless)	0.6	0.8	1
Inlet water temperature (K)	287.15	291.15	295.15
Water velocity (m·s ⁻¹)	1.5	1.8	2.1

be constant. Three values are chosen for each of the parameters: the steam mass flow rate, the noncondensable gases outlet pressure, the tube cleanliness factor, the inlet

water temperature, and the water velocity. These values give a sufficiently wide range of the condenser operating conditions.

3. ANALYSIS OF THE MODELS

In the present study, the methods are applied to a condenser for which the total heat transfer area, the geometry, the arrangement, the material and the geometric parameters of the tubes are known.

For the application of the SSSC to each set of operating conditions, the values of the following variables are required: the tube cleanliness factor, the water mass flow rate, which is calculated from the values of $T_{w,i}$ and u_w , and the heat transfer rate in the condenser, the value of which has been assumed to be equal to the value obtained using the numerical model for the comparison of the results of both methods. Equations (1)–(4) are solved and the condenser pressure, the overall heat transfer coefficient, the condensate mass flow rate, and the outlet water temperature are obtained.

For the application of the numerical model to each set of operating conditions, the values of the boundary conditions and the tube fouling thermal resistance are required. By solving equations (5)–(10), the distributions of the mixture velocity, the condenser pressure, the air mass fraction, the water and mixture temperatures, the heat transfer rate per unit volume, and the condensate mass flow rate per unit volume are obtained. From these results, the overall heat transfer coefficient, the total heat transfer rate in the condenser, the total condensate mass flow rate, and the average outlet water temperature are calculated.

For the comparison of both methods, a tube fouling thermal resistance, $R_{f,NM}$, is defined for the numerical model in accordance with the value of F_C used in the SSSC. This thermal resistance is calculated from equation (11) as the difference between the fouling thermal resistance of a tube, the area of which is equal to the total heat transfer area of the condenser, for a known cleanliness factor, $R_f(F_C)$, and the fouling thermal resistance of this tube for a cleanliness factor equal to 1, $R_f(F_C = 1)$. Thus, the fouling thermal resistance is equal to zero for clean tubes ($F_C = 1$) and increases when the cleanliness factor decreases:

$$R_{f,NM} = R_f(F_C) - R_f(F_C = 1) \quad (11)$$

The fouling thermal resistance of the tube is calculated assuming that the heat transfer rate in the tube is equal

to the heat transfer rate in the condenser. The former is obtained through equation (12), and the latter is calculated solving equations (1)–(4) from the values of m_c , F_C , $T_{w,i}$ and u_w , assuming that the total steam mass flow rate condenses. The value of m_c is equal to the steam mass flow rate used in the numerical model:

$$Q = \frac{1}{\frac{D_o}{D_i} \frac{1}{h_w} + \frac{D_o}{2k} \ln\left(\frac{D_o}{D_i}\right) + \frac{1}{h_c} + R_f(F_C)} A(T_g - T_{w,avg})$$

$$T_{w,avg} = \frac{T_{w,i} + T_{w,o}}{2} \quad (12)$$

With regard to the numerical model results, the total heat transfer rate in the condenser and the total condensate mass flow rate are calculated as the sum of the values obtained in each control volume from equations (13) and (14). The average value of the outlet water temperature is calculated from the average value of the water outlet specific enthalpy defined by equation (15):

$$Q = \frac{\sum_{j=1}^N q_{v,j} V_j}{\sum_{j=1}^N V_j} \quad (13)$$

$$m_c = \frac{\sum_{j=1}^N m_{c,v,j} V_j}{\sum_{j=1}^N V_j} \quad (14)$$

where $q_{v,j}$ is the heat rate per unit volume transferred in the control volume j , m_c is the condensate mass flow rate, $m_{c,v,j}$ is the condensate mass flow rate per unit volume in the control volume j , and V_j is the volume of the control volume j ;

$$i_{w,o,avg} = \frac{\sum_{j=1}^N i_{w,j} m_{w,j}}{\sum_{j=1}^N m_{w,j}} = \int_{273.15}^{T_{w,o,avg}} C_{p_w}(T) dT \quad (15)$$

where $i_{w,o}$ is the average outlet water specific enthalpy, $i_{w,j}$ is the average water specific enthalpy in the control volume j , $m_{w,j}$ is the water mass flow rate in the control volume j , and $T_{w,o,avg}$ is the average outlet water temperature.

With respect to the condenser pressure distribution, the SSSC assume a single value for each set of operating conditions, defined as the pressure maintained within the condenser shell at locations no more than 0.3 m from the first tube (P_{SSSC}). Figure 2 shows the pressure distribution obtained using the numerical model in two cross sections of the condenser, one situated near the water inlet, figure 2(a), and the other placed near the water outlet, figure 2(b). It can be observed that the pressure is maximum at the mixture inlet, and decreases as the mixture comes across the tubes to the venting outlet, where the pressure is minimum. Moreover, the

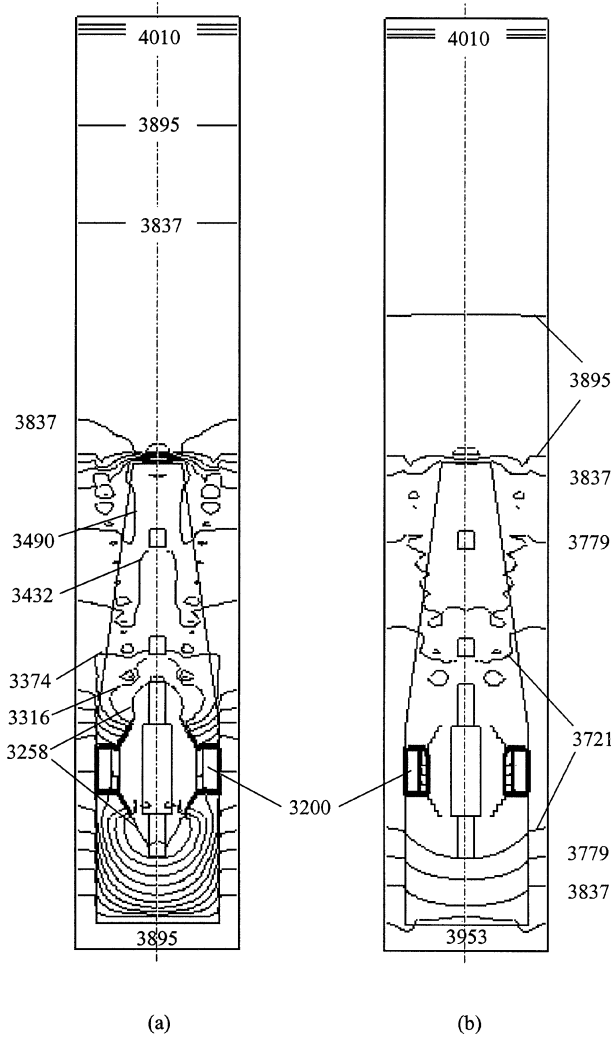


Figure 2. Distribution of the condenser pressure obtained with the numerical model (Pa): (a) XY plane at water inlet, (b) XY plane at water outlet.

pressure increases along the tubes as the water heats up, being higher at the water outlet than at the water inlet. Therefore, it is essential to define a criterion for the condenser pressure of the numerical model so as to be able to compare both methods, as the SSSC do not define the measurement point exactly. In the present study, the pressure in the control volume located in the cross section of the tube bundle situated at the middle of the tube length, on the condenser symmetry axis, and 0.05 m above the first tube row, is defined as the condenser pressure of the numerical model (P_{NM}), according to the SSSC.

The overall heat transfer coefficient of the numerical model (U_{NM}) is calculated from equation (3) from the

saturation temperature at pressure P_{NM} , and is compared with the overall heat transfer coefficient obtained using the SSSC, U_{SSSC} .

4. RESULTS

4.1. Condenser pressure

Figure 3 shows the variation of P_{SSSC} and P_{NM} versus u_w for $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$ and $P_{g,o} = 3200 \text{ Pa}$, varying the values of F_C and $T_{w,i}$. It is observed that P_{NM} is greater than P_{SSSC} in every case, both pressures decreasing when u_w and F_C increase, and when $T_{w,i}$ decreases.

In most cases, the difference in pressure between both methods presents little variation with respect to the parameters F_C , $T_{w,i}$ and u_w . However, in the cases corresponding to low values of $T_{w,i}$ and high values of F_C , P_{NM} is observed to behave differently, due to the value of $P_{g,o}$ chosen for the numerical simulations. Figure 4 shows the variation curves of P_{SSSC} and P_{NM} versus u_w for $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$, $F_C = 1.0$ and $T_{w,i} = 14^\circ\text{C}$, and for six different values of $P_{g,o}$. While the calculation method of P_{SSSC} does not depend on $P_{g,o}$, this parameter does affect the value of P_{NM} . Decreasing the value of $P_{g,o}$, the variation curves of P_{SSSC} and P_{NM} have a similar slope, and the difference between the pressures obtained using both methods decreases. Moreover, there is a value of $P_{g,o}$ (around 2600 Pa in this case) below which the pressure of the numerical model is considered independent of the said parameter. It was observed in later simulations that this value of $P_{g,o}$ increases with the value of $T_{w,i}$.

With regard to the variation of the condenser pressure with the value of the steam mass flow rate, table II presents the results obtained using both methods for $T_{w,i} = 22^\circ\text{C}$, $u_w = 1.8 \text{ m}\cdot\text{s}^{-1}$, and three different values of m_s , $P_{g,o}$ and F_C , as well as the relative error of pressure of the numerical model with respect to the SSSC. It can be seen that the relative error and the values of P_{SSSC} and P_{NM} decrease with the value of the steam mass flow rate, since the heat transfer rate and the cooling water mass flow rate are constant in each case.

4.2. Overall heat transfer coefficient

Figure 5 shows the variation of U_{SSSC} and U_{NM} versus u_w for $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$ and $P_{g,o} = 3200 \text{ Pa}$, varying

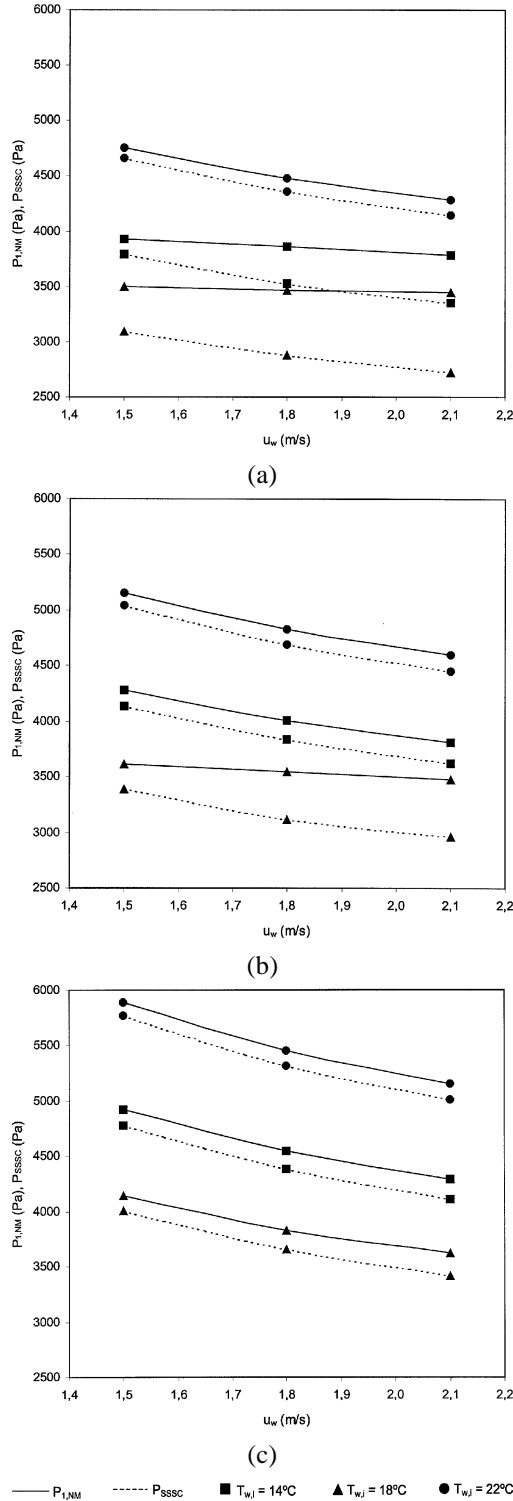


Figure 3. Variation of the condenser pressure with $P_{g,o}$, $T_{w,i}$ and u_w : $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$, (a) $F_C = 1.0$, (b) $F_C = 0.8$, and (c) $F_C = 0.6$.

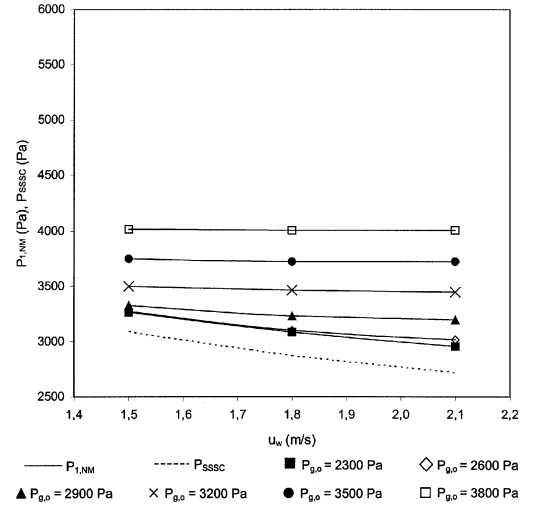


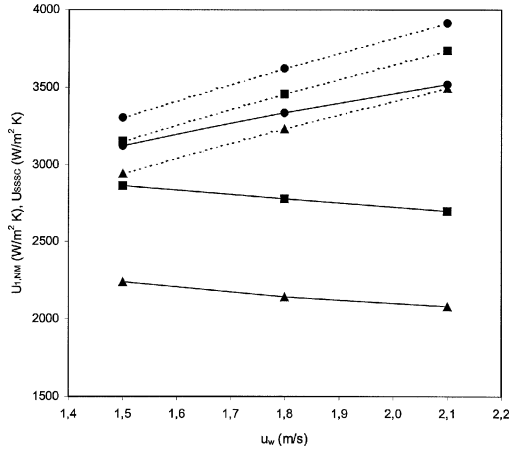
Figure 4. Variation of the condenser pressure with $P_{g,o}$ and u_w : $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$, $F_C = 1.0$, $T_{w,i} = 14^\circ\text{C}$.

TABLE II
Comparison of the condenser pressure ($T_{w,i} = 22^\circ\text{C}$, $u_w = 1.35 \text{ m}\cdot\text{s}^{-1}$).

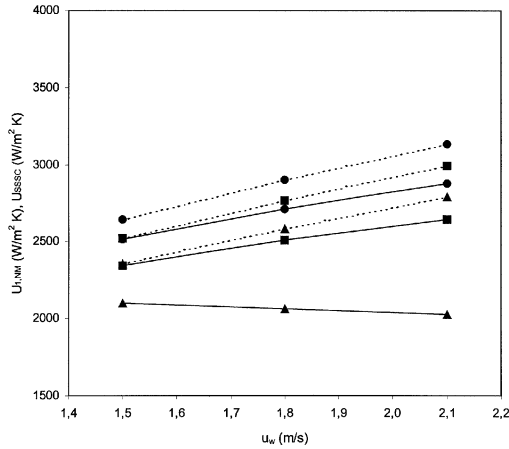
m_s ($\text{kg}\cdot\text{s}^{-1}$)	$P_{g,o}$ (Pa)	F_C	Condenser pressure (Pa)		Relative error (%)
			P_{NM}	$P_{SS,SC}$	ε_P
30.84	3800	1.0	4970	4658	2.70
30.84	3800	0.8	5371	5001	2.72
30.84	3800	0.6	6087	5642	2.50
23.73	3200	1.0	4331	4100	2.16
23.73	3200	0.8	4607	4335	2.37
23.73	3200	0.6	5091	4769	2.39
16.95	2800	1.0	3775	3621	1.36
16.95	2800	0.8	3952	3772	1.66
16.95	2800	0.6	4258	4046	1.86

the values of F_C and $T_{w,i}$. It is observed that U_{NM} is lower than $U_{SS,SC}$ in every case, both coefficients increasing when F_C increases and $T_{w,i}$ decreases.

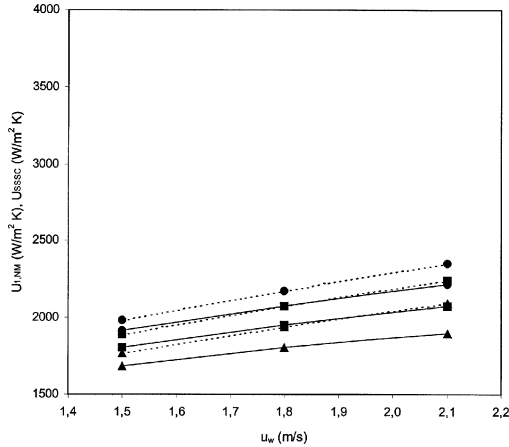
With respect to the variation of the overall heat transfer coefficient with u_w , both coefficients increase, in the majority of cases, with the value of u_w , the slope of the curve of $U_{SS,SC}$ being generally greater than that of U_{NM} . However, in the same cases as those observed in the pressure analysis, U_{NM} decreases for high values of u_w . This unexpected behaviour of U_{NM} is caused by the value of $P_{g,o}$ chosen for the numerical simulations. The influence of $P_{g,o}$ on the overall heat transfer coefficient was studied in a similar way to that of the pressure analysis. Figure 6 shows the variation curves of $U_{SS,SC}$ and U_{NM} versus u_w for $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$,



(a)



(b)



(c)

— $U_{1,NM}$ - - - - U_{SSSC} ■ $T_{w,i} = 14^\circ\text{C}$ ▲ $T_{w,i} = 18^\circ\text{C}$ ● $T_{w,i} = 22^\circ\text{C}$

Figure 5. Variation of the overall heat transfer coefficient with $P_{g,o}$, $T_{w,i}$ and u_w : $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$, (a) $F_C = 1.0$, (b) $F_C = 0.8$, and (c) $F_C = 0.6$.

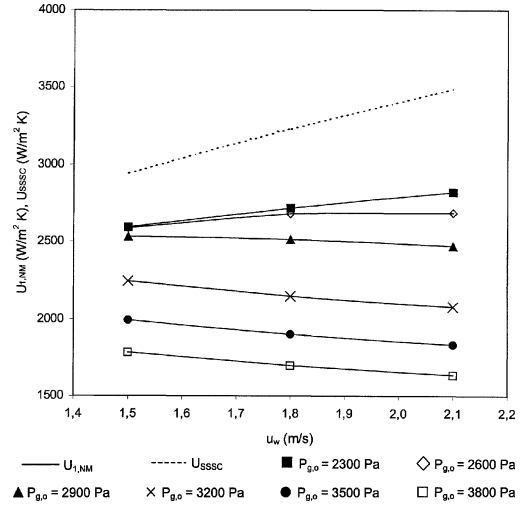


Figure 6. Variation of the overall heat transfer coefficient with $P_{g,o}$ and u_w : $m_s = 30.84 \text{ kg}\cdot\text{s}^{-1}$, $F_C = 1.0$, $T_{w,i} = 14^\circ\text{C}$.

TABLE III
Comparison of the overall heat transfer coefficient
($T_{w,i} = 22^\circ\text{C}$, $u_w = 1.35 \text{ m}\cdot\text{s}^{-1}$).

Parameters			Overall heat transfer coefficient ($\text{W}\cdot\text{m}^{-2}\cdot\text{K}^{-1}$)		Relative error (%)
m_s ($\text{kg}\cdot\text{s}^{-1}$)	$P_{g,o}$ (Pa)	F_C	U_{NM}	U_{SSSC}	ε_U
30.84	3800	1.0	2960.1	3615.1	8.0
30.84	3800	0.8	2429.1	2892.1	6.5
30.84	3800	0.6	1885.4	2169.1	4.6
23.73	3200	1.0	2918.3	3615.1	8.2
23.73	3200	0.8	2396.5	2892.1	7.2
23.73	3200	0.6	1863.1	2169.1	5.5
16.95	2800	1.0	2894.6	3615.1	7.1
16.95	2800	0.8	2375.1	2892.1	6.9
16.95	2800	0.6	1845.5	2169.1	5.9

$F_C = 1.0$ and $T_{w,i} = 14^\circ\text{C}$, and for six different values of $P_{g,o}$. While the calculation method of U_{SSSC} does not depend on $P_{g,o}$, this parameter does affect the value of U_{NM} . Decreasing the value of $P_{g,o}$, the slope of the curves of U_{NM} increases and its sign changes from negative to positive, decreasing the differences between both methods. Moreover, there is a value of $P_{g,o}$ (the same as observed in the pressure analysis) below which the overall heat transfer coefficient of the numerical model is considered independent of the said parameter.

In relation to the variation of the overall heat transfer coefficient with the value of the steam mass flow rate, *table III* presents the results obtained using both methods for $T_{w,i} = 22^\circ\text{C}$, $u_w = 1.8 \text{ m}\cdot\text{s}^{-1}$, and three different values of m_s , $P_{g,o}$ and F_C , as well as the relative error

of the numerical model with respect to the SSSC. It can be seen that the values of U_{SSC} and U_{NM} decrease with the value of the steam mass flow rate, while the relative error is almost constant.

5. CONCLUSIONS

In the present study, the variables *condenser pressure* and *overall heat transfer coefficient* obtained using the "Standards for Steam Surface Condensers" (SSSC) of the Heat Exchange Institute and a three-dimensional numerical model were compared. The comparison was carried out from two different points of view: the evolution of the variables with five parameters (the steam mass flow rate from the turbine exhaust, the venting outlet pressure, the tube cleanliness factor, the inlet water temperature, and the water velocity), as well as the quantification of the differences between the values obtained using both methods.

As expected, both in the numerical model and in the SSSC, the condenser pressure decreases and the overall heat transfer coefficient increases when the heat transfer processes in the condenser are improved due to the variation of certain parameters. In particular, this tendency appears when there is an increment in the water velocity, and when the tube cleanliness factor or the inlet water temperature decreases. With respect to the steam mass flow rate, the condenser pressure and the overall heat transfer coefficient obtained using both methods decrease with this parameter.

Tendencies differing from this overall performance, which are observed in the numerical results for low values of the inlet water temperature and high values of the tube cleanliness factor, are corrected by decreasing the value of the venting outlet pressure.

In spite of the similar overall behaviour of the methods, there are some differences between the values of the condenser pressure and the overall heat transfer coefficient obtained using both methods. The numerical model predicts higher condenser pressures and lower overall heat transfer coefficients than the SSSC. Consequently, the analysis of the condenser performance will be more conservative using the numerical model than that of the SSSC, and the design of condensers with the numerical model will lead to larger equipment than that obtained using the SSSC method, a conclusion that is contrast with the logical thought that simplified engineering calculation methods adopt a certain degree of conservatism, being used only for gross estimates. There are two main reasons for this conclusion: on the one hand,

the SSSC do not take into account certain negative effects that aggravate the heat and mass transfer in the condensers, mainly the inundation of the tubes by the condensate and the presence of noncondensable gases mixed with the steam. On the other hand, the details concerning the performance conditions and the geometries to which the SSSC could be satisfactorily applied are very poor.

With regards to the steam mass flow rate, the differences between the condenser pressures obtained using both methods decrease with this parameter, while the differences between the overall heat transfer coefficients are almost constant. The reason is that when the steam mass flow rate decreases and the values of the heat transfer rate and the cooling water mass flow rate are maintained constant, the pressure predicted by the numerical model decreases, approaching the pressure predicted by that of the SSSC, but satisfactorily predicting the overall heat transfer coefficient.

The conclusions of this study could be extended to similar numerical models, although it may be possible that numerical models applied to condensers operating under performance conditions different from those simulated in this study, might lead to some variations in the overall behaviour.

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